

1. (a) $N_{1a}^c = \delta_{ac}$

This property is the trivial statement that particles do not undergo fusion with the vacuum. This is the fundamental property of the vacuum. The δ function states that the only outcome of fusing the particle with the vacuum is to obtain the particle itself, which is the same as having no fusion take place.

(b) $N_{ab}^1 = \delta_{b\bar{a}}$

This property states that two particles fuse to produce vacuum (annihilate), only if they are antiparticles. If particle b annihilates to the vacuum with particle a . Then by definition, this particle is the antiparticle of a , so it can be referred to as $b=\bar{a}$. A particle/antiparticle pair, have the vacuum as their unique fusion channel, thus giving rise to the δ relation.

(c) $N_{ab}^c = N_{ba}^c = N_{b\bar{c}}^{\bar{a}} = N_{\bar{a}\bar{b}}^{\bar{c}}$

i. $N_{ab}^c = N_{ba}^c$

The first equality here follows directly from the commutativity property of the fusion rules, specifically that $a \times b = b \times a$. Therefore by changing the order of fusion between two particles, the fusion channel should remain identical.

ii. $N_{ab}^c = N_{b\bar{c}}^{\bar{a}}$

$$\begin{aligned} (a \times b) \times (\bar{c} \times \bar{a}) &= \left(\sum_j N_{ab}^j j \right) \times (\bar{c} \times \bar{a}) = \sum_j N_{ab}^j (j \times \bar{c} \times \bar{a}) \\ &= N_{ab}^c (c \times \bar{c}) \times \bar{a} + \sum_{j \neq c} N_{ab}^j (j \times \bar{c} \times \bar{a}) \\ &= N_{ab}^c \bar{a} + \sum_{j \neq c} N_{ab}^j (j \times \bar{c} \times \bar{a}) \end{aligned}$$

$$\begin{aligned} (a \times b) \times (\bar{c} \times \bar{a}) &= (a \times \bar{a}) \times (b \times \bar{c}) \\ &= (1) \times \left(\sum_k N_{b\bar{c}}^k k \right) \\ &= N_{b\bar{c}}^{\bar{a}} \bar{a} + \sum_{k \neq \bar{a}} N_{b\bar{c}}^k k \\ &= N_{ab}^c \bar{a} + \sum_{j \neq c} N_{ba}^j (j \times \bar{c} \times \bar{a}) \end{aligned}$$

We would like here to equate the coefficients of the leading \bar{a} terms, however we must first eliminate the possibility of the final fusion term resulting in a \bar{a} anyon. $\sum_{j \neq c} N_{ab}^j (j \times \bar{c} \times \bar{a})$.

By the associativity of the fusion rules, $j \times \bar{c} \times \bar{a} = (j \times \bar{c}) \times \bar{a}$. Due to the fact that anyons do not fuse with the vacuum and change to a different anyon, we have $N_{mk}^k = \delta_{m1}$. Thus for $(j \times \bar{c}) \times \bar{a}$ to result in a \bar{a} anyon, it must be that $j \times \bar{c} = 1$. As an anyon/anti-anyon pair have a unique fusion channel with the vacuum we have $N_m^1 k = \delta_{m\bar{k}}$. Therefore in order for $j \times \bar{c} = 1$, it must be that $j = c$. This is seen to be the original term isolated from the summation, therefore for $j \neq c$, we can not have $j \times \bar{c} \times \bar{a}$ resulting in a \bar{a} anyon.

Finally by equating the coefficients of \bar{a} it is seen that $N_{ab}^c = N_{b\bar{c}}^{\bar{a}}$. As required.

iii. $N_{ab}^c = N_{b\bar{c}}^{\bar{a}}$

In order to demonstrate this equality we must show that some intermediate results hold. The process will be as follows:

$$N_{ab}^c = N_{a\bar{c}}^{\bar{b}} = N_{ab\bar{c}}^1 = N_a^{\bar{b}c} = N_c^{\bar{a}\bar{b}} = N_{\bar{a}\bar{b}}^{\bar{c}}$$

Firstly we have

$$N_{ab}^c = N_{a\bar{c}}^{\bar{b}}$$

from the earlier equality.

The second equality defines what it means to have three anyons undergoing fusion. To show this equality consider:

$$a \times b \times c = \sum_i \sum_j N_{ab}^i N_{ic}^j = N_{abc}^A$$

Where A is the outcome of fusing all three anyons. In this case we are interested in $A = 1$, the vacuum. We split the above summation to the case $j = 1$ and $j \neq 1$. In the former case, we are forced to take $i = c$, due to the delta relation:

$$N_{ab}^1 = \delta_{b\bar{a}}$$

which dictates that the only way to have the vacuum as outcome from a fusion is to fuse two antiparticles. We then have:

$$N_{ab}^c + \sum_{i \neq c} \sum_{j \neq 1} N_{ab}^i N_{ic}^j$$

The first term here was obtained when we considered fusing three anyons a, b, c to give the vacuum, therefore it can be expressed as $N_{ab\bar{c}}^1$, and we have $N_{ab}^c = N_{ab\bar{c}}^1$. Next, we consider the composite anyon $b \times \bar{c}$. We are able to manipulate this without evaluating the exact outcome of the fusion. We have that

$$b \times c \times \bar{b} \times \bar{c} = b \times \bar{b} \times c \times \bar{c} = (b \times \bar{b}) \times (c \times \bar{c}) = 1$$

by the commutativity and associativity properties of fusion. Therefore we know the antiparticle of the $b \times \bar{c}$ anyon is simply $\bar{b} \times c$.

Therefore by the earlier property (ii), we are able to raise the index: $b\bar{c}$, by taking its antiparticle and lowering the 1 anyon (the vacuum) which is its own antiparticle. This gives

$$N_{ab\bar{c}}^1 = N_a^{\bar{b}c}$$

By applying this property again, to the a, c anyons, we obtain

$$N_a^{\bar{b}c} = N_{\bar{c}}^{\bar{b}a}$$

The final step is to recall that the fusion process is time invertible as $a \times b = c$ can equivalently be seen as the particle c splitting to give a, b . Therefore we have that:

$$N_{\bar{a}}^{\bar{b}\bar{c}} = N_{\bar{b}\bar{c}}^{\bar{a}}$$

by time inverting the process. And consequently by considering the string of equalities demonstrated:

$$N_{ab}^c = N_{a\bar{c}}^{\bar{b}} = N_{ab\bar{c}}^1 = N_a^{\bar{b}c} = N_{\bar{c}}^{\bar{b}a} = N_{\bar{a}\bar{b}}^{\bar{c}}$$

we have that:

$$N_{ab}^c = N_{\bar{b}\bar{c}}^{\bar{a}}$$

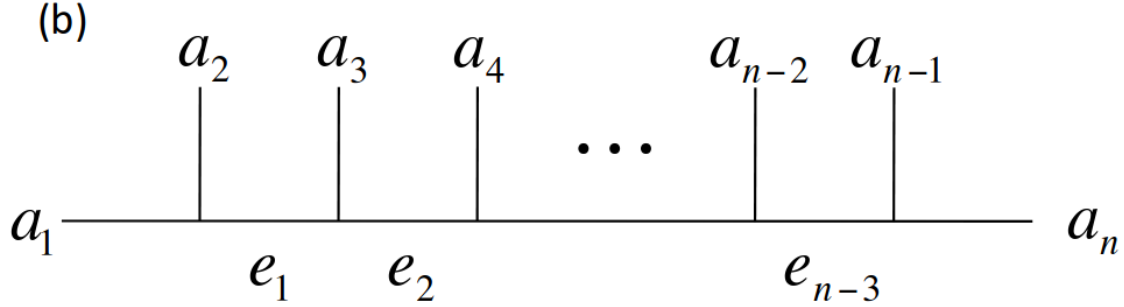
(d)

$$(a \times b) \times c = \left(\sum_d N_{ab}^d e \right) \times c = \sum_d \left(\sum_e N_{ab}^e N_{ec}^d \right) d$$

$$\begin{aligned} a \times (b \times c) &= a \times \sum_f N_{bc}^f f = \sum_d N_{bc}^d (a \times f) \\ &= \sum_d \left(\sum_f N_{bc}^f N_{af}^d \right) d = \sum_d \left(\sum_f N_{af}^d N_{bc}^f \right) d \end{aligned}$$

These fusion rules are associative, meaning $(a \times b) \times c = a \times (b \times c)$, by equating the bracketed coefficients of d in each case it is seen that $\sum_e N_{ab}^e N_{ec}^d = \sum_f N_{bc}^f N_{af}^d$ as required.

2. Consider the fusion of n anyons.



The Hilbert space defining the fusion process will therefore be:

$$\mathcal{H}_{a_1, a_2, \dots, a_{n-1}}^{e_n} = \bigoplus_{e_1, e_2, \dots, e_n} \mathcal{H}_{a_1 a_2}^{e_1} \otimes \mathcal{H}_{e_1 a_3}^{e_2} \otimes \mathcal{H}_{e_2 a_4}^{e_3} \otimes \dots \otimes \mathcal{H}_{a_{n-1} e_{n-3}}^{a_n}$$

This new Hilbert space will have dimension equal to:

$$\dim(\mathcal{M}_n) = \sum_{e_i} N_{a_1 a_2}^{e_1} N_{e_1 a_3}^{e_2} N_{e_2 a_4}^{e_3} \dots N_{a_{n-1} e_{n-3}}^{a_n}$$

If we consider the matrix N_c with entries $(N_c)_{ab} = N_{ab}^c$, we see that:

$$\dim(\mathcal{M}_n) = \sum_{e_i} N_{a_1 a_2}^{e_1} N_{e_1 a_3}^{e_2} N_{e_2 a_4}^{e_3} \dots N_{a_{n-1} e_{n-3}}^{a_n} = (N_{a_2} N_{a_3} \dots N_{a_n})_{a_1}^{b_n}$$

We then have

$$\dim(\mathcal{H}_{a_1, a_2, \dots, a_{n-1}}^{e_n}) = (N_{a_2} N_{a_3} \dots N_{a_n})_{a_1}^{e_n} = [N_a^{(n-1)}]_a^{e_n} \approx d_a^n$$

3. It is possible to create an entangled state, however it is **not** possible through topologically protected operations alone. More information can be found here:

Brennen, G. K., Iblisdir, S., Pachos, J. K. and Slingerland, J. K. 2009. New J. Phys. 11, 103023

4. Pentagon Equation

F matrix for the Fibonacci model:

$$F_{\tau\tau\tau}^\tau = \begin{pmatrix} \frac{1}{\phi} & \frac{1}{\sqrt{\phi}} \\ \frac{1}{\sqrt{\phi}} & -\frac{1}{\phi} \end{pmatrix} \quad (1)$$

Pentagon Equation:

$$(F_{12c}^5)_a^d (F_{a34}^5)_b^c = \sum_e (F_{234}^d)_e^c (F_{1e4}^5)_b^d (F_{123}^b)_a^e \quad (2)$$

The only non-trivial pentagon arises when $1 = 2 = 3 = 4 = 5 = \tau$ [1]. The pentagon equation (2) therefore becomes:

$$(F_{\tau\tau c}^\tau)_a^d (F_{a\tau\tau}^\tau)_b^c = \sum_e (F_{\tau\tau\tau}^d)_e^c (F_{\tau e\tau}^\tau)_b^d (F_{\tau\tau\tau}^b)_a^e$$

Let $a = 1$.

$$(F_{\tau\tau c}^\tau)_1^d (F_{1\tau\tau}^\tau)_b^c = (F_{\tau\tau\tau}^d)_1^c (F_{\tau 1\tau}^\tau)_b^d (F_{\tau\tau\tau}^b)_1^1 + (F_{\tau\tau\tau}^d)_\tau^c (F_{\tau\tau\tau}^\tau)_b^d (F_{\tau\tau\tau}^b)_1^\tau$$

Let $b = 1$.

$$(F_{\tau\tau c}^\tau)_1^d (F_{1\tau\tau}^\tau)_1^c = (F_{\tau\tau\tau}^d)_1^c (F_{\tau 1\tau}^\tau)_1^d (F_{\tau\tau\tau}^1)_1^1 + (F_{\tau\tau\tau}^d)_\tau^c (F_{\tau\tau\tau}^\tau)_1^d (F_{\tau\tau\tau}^1)_1^\tau$$

Here we can identify that $(F_{1\tau\tau}^\tau)_1^c = (F_{\tau\tau\tau}^1)_1^\tau = (F_{\tau 1\tau}^\tau)_1^d = 0$ as they each correspond to an impossible fusion process. The pentagon equation therefore reduces to $0 = 0$ and we move on.

Let $a = 1, b = \tau$.

$$(F_{\tau\tau c}^\tau)_1^d (F_{1\tau\tau}^\tau)_\tau^c = (F_{\tau\tau\tau}^d)_1^c (F_{\tau 1\tau}^\tau)_\tau^d (F_{\tau\tau\tau}^\tau)_1^1 + (F_{\tau\tau\tau}^d)_\tau^c (F_{\tau\tau\tau}^\tau)_\tau^d (F_{\tau\tau\tau}^\tau)_1^\tau$$

Let $a = 1, b = \tau, c = 1$

$$(F_{\tau\tau 1}^\tau)_1^d (F_{1\tau\tau}^\tau)_\tau^1 = (F_{\tau\tau\tau}^d)_1^1 (F_{\tau 1\tau}^\tau)_\tau^d (F_{\tau\tau\tau}^\tau)_1^1 + (F_{\tau\tau\tau}^d)_\tau^1 (F_{\tau\tau\tau}^\tau)_\tau^d (F_{\tau\tau\tau}^\tau)_1^\tau$$

Here $(F_{\tau\tau 1}^\tau)_1^d = (F_{1\tau\tau}^\tau)_\tau^1 = 0$ as they each corresponds to an impossible fusion process. So we have:

$$0 = (F_{\tau\tau\tau}^d)_1^1 (F_{\tau 1\tau}^\tau)_\tau^d (F_{\tau\tau\tau}^\tau)_1^1 + (F_{\tau\tau\tau}^d)_\tau^1 (F_{\tau\tau\tau}^\tau)_\tau^d (F_{\tau\tau\tau}^\tau)_1^\tau$$

Which means:

$$(F_{\tau\tau\tau}^d)_1^1 (F_{\tau 1\tau}^\tau)_\tau^d (F_{\tau\tau\tau}^\tau)_1^1 = -(F_{\tau\tau\tau}^d)_\tau^1 (F_{\tau\tau\tau}^\tau)_\tau^d (F_{\tau\tau\tau}^\tau)_1^\tau$$

Let $a = 1, b = \tau, c = 1, d = 1$.

$$(F_{\tau\tau\tau}^1)_1^1 (F_{\tau 1\tau}^\tau)_\tau^1 (F_{\tau\tau\tau}^\tau)_1^1 = -(F_{\tau\tau\tau}^1)_\tau^1 (F_{\tau\tau\tau}^\tau)_\tau^1 (F_{\tau\tau\tau}^\tau)_1^\tau$$

Here we have $(F_{\tau 1\tau}^\tau)_\tau^1 = (F_{\tau\tau\tau}^\tau)_\tau^1 = 0$ as they correspond to impossible fusion processes. The pentagon equation therefore reduces to $0 = 0$ and we move on.

Let $a = 1, b = \tau, c = 1, d = \tau$.

$$(F_{\tau\tau\tau}^\tau)_1^1 (F_{\tau 1\tau}^\tau)_\tau^\tau (F_{\tau\tau\tau}^\tau)_1^1 = -(F_{\tau\tau\tau}^\tau)_\tau^1 (F_{\tau\tau\tau}^\tau)_\tau^\tau (F_{\tau\tau\tau}^\tau)_1^\tau$$

We can see here that $(F_{\tau 1\tau}^\tau)_\tau^\tau = 1$ by the fusion diagrams. Upon substitution we arrive at an equation relating elements of the $F_{\tau\tau\tau}^\tau$ matrix:

$$(F_{\tau\tau\tau}^\tau)_1^1 (F_{\tau\tau\tau}^\tau)_1^1 = -(F_{\tau\tau\tau}^\tau)_\tau^1 (F_{\tau\tau\tau}^\tau)_\tau^\tau (F_{\tau\tau\tau}^\tau)_1^\tau \quad (3)$$

Letting $a = \tau$, the pentagon equation (2) becomes:

$$(F_{\tau\tau c}^\tau)_\tau^d (F_{\tau\tau\tau}^\tau)_b^c = (F_{\tau\tau\tau}^d)_1^c (F_{\tau 1\tau}^\tau)_b^d (F_{\tau\tau\tau}^b)_\tau^1 + (F_{\tau\tau\tau}^d)_\tau^c (F_{\tau\tau\tau}^\tau)_b^d (F_{\tau\tau\tau}^b)_\tau^\tau$$

Let $a = \tau, b = 1$

$$(F_{\tau\tau c}^\tau)_\tau^d (F_{\tau\tau\tau}^\tau)_1^c = (F_{\tau\tau\tau}^d)_1^c (F_{\tau 1\tau}^\tau)_1^d (F_{\tau\tau\tau}^1)_\tau^1 + (F_{\tau\tau\tau}^d)_\tau^c (F_{\tau\tau\tau}^\tau)_1^d (F_{\tau\tau\tau}^1)_\tau^\tau$$

We have $(F_{\tau 1\tau}^\tau)_1^d = 0$ as it corresponds to an impossible fusion process, also $(F_{\tau\tau\tau}^1)_\tau^\tau = 1$ by the fusion diagram. Therefore:

$$(F_{\tau\tau c}^\tau)_\tau^d (F_{\tau\tau\tau}^\tau)_1^c = (F_{\tau\tau\tau}^d)_\tau^c (F_{\tau\tau\tau}^\tau)_1^d$$

Let $a = \tau, b = 1, c = 1$

$$(F_{\tau\tau 1}^\tau)^\tau (F_{\tau\tau\tau}^\tau)_1^1 = (F_{\tau\tau\tau}^d)_\tau^1 (F_{\tau\tau\tau}^\tau)_1^d$$

Let $a = \tau, b = 1, c = 1, d = 1$

$$(F_{\tau\tau 1}^\tau)_\tau^1 (F_{\tau\tau\tau}^\tau)_1^1 = (F_{\tau\tau\tau}^1)_\tau^1 (F_{\tau\tau\tau}^\tau)_1^1$$

Here we see that $(F_{\tau\tau 1}^\tau)_\tau^1 = (F_{\tau\tau\tau}^1)_\tau^1 = 0$ as they correspond to impossible fusion processes. Thus the pentagon equation reduces to $0=0$ and we move on.

Let $a = \tau, b = 1, c = 1, d = \tau$

$$(F_{\tau\tau 1}^\tau)_\tau^\tau (F_{\tau\tau\tau}^\tau)_1^1 = (F_{\tau\tau\tau}^\tau)_\tau^1 (F_{\tau\tau\tau}^\tau)_1^\tau$$

We can identify here that $(F_{\tau\tau 1}^\tau)_\tau^\tau = 1$ by the fusion diagrams, therefore we obtain another equation relating elements of the $F_{\tau\tau\tau}^\tau$ matrix.

$$(F_{\tau\tau\tau}^\tau)_1^1 = (F_{\tau\tau\tau}^\tau)_\tau^1 (F_{\tau\tau\tau}^\tau)_1^\tau \quad (4)$$

Consider $a = \tau, b = 1, c = \tau$

$$(F_{\tau\tau\tau}^\tau)_\tau^d (F_{\tau\tau\tau}^\tau)_1^\tau = (F_{\tau\tau\tau}^d)_\tau^\tau (F_{\tau\tau\tau}^\tau)_1^d$$

Let $d = 1$

$$(F_{\tau\tau\tau}^\tau)_\tau^1 (F_{\tau\tau\tau}^\tau)_1^\tau = (F_{\tau\tau\tau}^1)_\tau^\tau (F_{\tau\tau\tau}^\tau)_1^1$$

Identifying $(F_{\tau\tau\tau}^1)_\tau^\tau = 1$ we see that we have reached the same equation as above (4).
Let $d = \tau$

$$(F_{\tau\tau\tau}^\tau)_\tau^\tau (F_{\tau\tau\tau}^\tau)_1^\tau = (F_{\tau\tau\tau}^\tau)_\tau^\tau (F_{\tau\tau\tau}^\tau)_1^\tau$$

This is a trivial equality and of no further use.

Let $a = \tau, b = \tau, c = 1$

$$(F_{\tau\tau 1}^\tau)_\tau^d (F_{\tau\tau\tau}^\tau)_\tau^1 = (F_{\tau\tau\tau}^d)_\tau^1 (F_{\tau 1 \tau}^\tau)_\tau^d (F_{\tau\tau\tau}^\tau)_\tau^1 + (F_{\tau\tau\tau}^d)_\tau^1 (F_{\tau\tau\tau}^\tau)_\tau^d (F_{\tau\tau\tau}^\tau)_\tau^1$$

Let $d = 1$

$$(F_{\tau\tau 1}^\tau)_\tau^1 (F_{\tau\tau\tau}^\tau)_\tau^1 = (F_{\tau\tau\tau}^1)_\tau^1 (F_{\tau 1 \tau}^\tau)_\tau^1 (F_{\tau\tau\tau}^\tau)_\tau^1 + (F_{\tau\tau\tau}^1)_\tau^1 (F_{\tau\tau\tau}^\tau)_\tau^1 (F_{\tau\tau\tau}^\tau)_\tau^1$$

We can see here that $(F_{\tau\tau 1}^\tau)_\tau^1 = (F_{\tau 1 \tau}^\tau)_\tau^1 = (F_{\tau\tau\tau}^1)_\tau^1 = 0$ as they each correspond to an impossible fusion process. The pentagon equation therefore reduces to $0=0$ and we move on.

Let $d = \tau$

$$(F_{\tau\tau 1}^\tau)_\tau^\tau (F_{\tau\tau\tau}^\tau)_\tau^1 = (F_{\tau\tau\tau}^\tau)_\tau^1 (F_{\tau 1 \tau}^\tau)_\tau^\tau (F_{\tau\tau\tau}^\tau)_\tau^1 + (F_{\tau\tau\tau}^\tau)_\tau^1 (F_{\tau\tau\tau}^\tau)_\tau^\tau (F_{\tau\tau\tau}^\tau)_\tau^1$$

Here we have $(F_{\tau\tau 1}^\tau)^\tau = (F_{\tau 1\tau}^\tau)^\tau = 1$ by the fusion diagrams. Upon substitution we arrive at a third equation relating elements of the $F_{\tau\tau\tau}^\tau$ matrix:

$$(F_{\tau\tau\tau}^\tau)_\tau^1 = (F_{\tau\tau\tau}^\tau)_1^1 (F_{\tau\tau\tau}^\tau)_\tau^1 + (F_{\tau\tau\tau}^\tau)_\tau^1 (F_{\tau\tau\tau}^\tau)_\tau^\tau (F_{\tau\tau\tau}^\tau)_\tau^\tau \quad (5)$$

Let $a = b = c = \tau, d = 1$

$$(F_{\tau\tau\tau}^\tau)_\tau^1 (F_{\tau\tau\tau}^\tau)_\tau^\tau = (F_{\tau\tau\tau}^\tau)_1^\tau (F_{\tau 1\tau}^\tau)_\tau^1 (F_{\tau\tau\tau}^\tau)_\tau^1 + (F_{\tau\tau\tau}^\tau)_\tau^1 (F_{\tau\tau\tau}^\tau)_\tau^\tau (F_{\tau\tau\tau}^\tau)_\tau^\tau$$

Here $(F_{\tau 1\tau}^\tau)_\tau^1 = 0$ and $(F_{\tau\tau\tau}^\tau)_\tau^\tau = 1$, giving:

$$(F_{\tau\tau\tau}^\tau)_\tau^1 (F_{\tau\tau\tau}^\tau)_\tau^\tau = (F_{\tau\tau\tau}^\tau)_\tau^1 (F_{\tau\tau\tau}^\tau)_\tau^\tau$$

which is trivial and of no further use.

The final case: $a = b = c = d = \tau$:

$$(F_{\tau\tau\tau}^\tau)_\tau^\tau (F_{\tau\tau\tau}^\tau)_\tau^\tau = (F_{\tau\tau\tau}^\tau)_1^\tau (F_{\tau 1\tau}^\tau)_\tau^\tau (F_{\tau\tau\tau}^\tau)_\tau^1 + (F_{\tau\tau\tau}^\tau)_\tau^\tau (F_{\tau\tau\tau}^\tau)_\tau^\tau (F_{\tau\tau\tau}^\tau)_\tau^\tau$$

Here $(F_{\tau 1\tau}^\tau)_\tau^\tau = 1$ by the fusion rules, giving one final equation relating the elements of the $F_{\tau\tau\tau}^\tau$ matrix:

$$(F_{\tau\tau\tau}^\tau)_\tau^\tau (F_{\tau\tau\tau}^\tau)_\tau^\tau = (F_{\tau\tau\tau}^\tau)_1^\tau (F_{\tau\tau\tau}^\tau)_\tau^1 + (F_{\tau\tau\tau}^\tau)_\tau^\tau (F_{\tau\tau\tau}^\tau)_\tau^\tau (F_{\tau\tau\tau}^\tau)_\tau^\tau \quad (6)$$

Let us now consider the given $F_{\tau\tau\tau}^\tau$ matrix for the Fibonacci model (1).

Substituting this into the first equation we derived (3), we see that:

$$\frac{1}{\phi} \cdot \frac{1}{\phi} = - \left(\frac{1}{\sqrt{\phi}} \cdot \frac{-1}{\phi} \cdot \frac{1}{\sqrt{\phi}} \right)$$

It is seen that this equation is satisfied for all ϕ .

Substituting (1) into the second equation we derived (4) gives:

$$\frac{1}{\phi} = \frac{1}{\sqrt{\phi}} \cdot \frac{1}{\sqrt{\phi}}$$

This equation is satisfied for all ϕ .

Substituting (1) into the third equation we derived (5) gives:

$$\frac{1}{\sqrt{\phi}} = \left(\frac{1}{\phi} \cdot \frac{1}{\sqrt{\phi}} \right) + \left(\frac{1}{\sqrt{\phi}} \cdot \frac{-1}{\phi} \cdot \frac{-1}{\phi} \right)$$

This is satisfied for $\phi = \frac{1+\sqrt{5}}{2}$: the golden ratio.

Substituting (1) into the final equation we derived (6) gives:

$$\frac{1}{\phi^2} = \frac{1}{\phi} + \left(\frac{-1}{\phi} \cdot \frac{-1}{\phi} \cdot \frac{-1}{\phi} \right)$$

This is satisfied for $\phi = \frac{1+\sqrt{5}}{2}$: the golden ratio.

Thus the given F matrix for the Fibonacci model satisfies the Pentagon equation.

Hexagon Equation

R matrix for the Fibonacci model:

$$R_{\tau\tau} = \begin{pmatrix} e^{4\pi i/5} & 0 \\ 0 & -e^{2\pi i/5} \end{pmatrix} \quad (7)$$

Hexagon Equation:

$$\sum_b (F_{231}^4)_b^c R_{1b}^4 (F_{123}^4)_a^b = R_{13}^c (F_{213}^4)_1^c R_{12}^a \quad (8)$$

As in the pentagon equation, set $1 = 2 = 3 = 4 = 5 = \tau$. The hexagon equation (8) becomes:

$$(F_{\tau\tau\tau}^\tau)_\tau^c R_{\tau\tau}^\tau (F_{\tau\tau\tau}^\tau)_a^\tau + (F_{\tau\tau\tau}^\tau)_1^c R_{\tau 1}^\tau (F_{\tau\tau\tau}^\tau)_a^1 = R_{\tau\tau}^c (F_{\tau\tau\tau}^\tau)_a^c R_{\tau\tau}^a \quad (9)$$

There are four separate cases for all pairs $(a, c) \in \{1, \tau\}^2$ which must be considered.

(a) $a = c = 1$

$$(F_{\tau\tau\tau}^\tau)_\tau^1 R_{\tau\tau}^\tau (F_{\tau\tau\tau}^\tau)_1^1 + (F_{\tau\tau\tau}^\tau)_1^1 R_{\tau 1}^\tau (F_{\tau\tau\tau}^\tau)_1^1 = R_{\tau\tau}^1 (F_{\tau\tau\tau}^\tau)_1^1 R_{\tau\tau}^1$$

It is possible to identify that $R_{\tau 1}^\tau = 1$.

Substituting from the $F_{\tau\tau\tau}^\tau$ matrix we obtain an equation relating the components of the $R_{\tau\tau}$ matrix.

$$\frac{1}{\phi^2} + \frac{1}{\phi} R_{\tau\tau}^\tau = \frac{1}{\phi} R_{\tau\tau}^1 \cdot R_{\tau\tau}^1 \quad (10)$$

(b) $a = \tau, c = 1$

$$(F_{\tau\tau\tau}^\tau)_\tau^1 R_{\tau\tau}^\tau (F_{\tau\tau\tau}^\tau)_\tau^\tau + (F_{\tau\tau\tau}^\tau)_1^1 R_{\tau 1}^\tau (F_{\tau\tau\tau}^\tau)_\tau^1 = R_{\tau\tau}^1 (F_{\tau\tau\tau}^\tau)_\tau^1 R_{\tau\tau}^\tau$$

It is possible to identify that $R_{\tau 1}^\tau = 1$.

Substituting from the $F_{\tau\tau\tau}^\tau$ matrix we obtain a second equation relating the components of the $R_{\tau\tau}$ matrix.

$$\frac{1}{\phi\sqrt{\phi}} - \frac{1}{\phi\sqrt{\phi}} R_{\tau\tau}^\tau = \frac{1}{\sqrt{\phi}} R_{\tau\tau}^1 R_{\tau\tau}^\tau \quad (11)$$

(c) $a = 1, c = \tau$

$$(F_{\tau\tau\tau}^\tau)_\tau^\tau R_{\tau\tau}^\tau (F_{\tau\tau\tau}^\tau)_1^\tau + (F_{\tau\tau\tau}^\tau)_1^\tau R_{\tau 1}^\tau (F_{\tau\tau\tau}^\tau)_1^1 = R_{\tau\tau}^\tau (F_{\tau\tau\tau}^\tau)_1^\tau R_{\tau\tau}^1$$

It is possible to identify that $R_{\tau 1}^\tau = 1$.

Substituting from the $F_{\tau\tau\tau}^\tau$ matrix we obtain a third equation relating the components of the $R_{\tau\tau}$ matrix.

$$\frac{1}{\phi\sqrt{\phi}} - \frac{1}{\phi\sqrt{\phi}} R_{\tau\tau}^\tau = \frac{1}{\sqrt{\phi}} R_{\tau\tau}^1 R_{\tau\tau}^\tau \quad (12)$$

This is equal to the second equation obtained (11).

(d) $a = \tau, c = \tau$

$$(F_{\tau\tau\tau}^\tau)^\tau R_{\tau\tau}^\tau (F_{\tau\tau\tau}^\tau)^\tau + (F_{\tau\tau\tau}^\tau)_1^\tau R_{\tau 1}^\tau (F_{\tau\tau\tau}^\tau)_1^\tau = R_{\tau\tau}^\tau (F_{\tau\tau\tau}^\tau)^\tau R_{\tau\tau}^\tau$$

It is possible to identify that $R_{\tau 1}^\tau = 1$.

Substituting from the $F_{\tau\tau\tau}^\tau$ matrix we obtain a fourth and final equation relating the components of the $R_{\tau\tau}^\tau$ matrix.

$$\frac{1}{\phi} + \frac{1}{\phi^2} R_{\tau\tau}^\tau = -\frac{1}{\phi} R_{\tau\tau}^\tau R_{\tau\tau}^\tau \quad (13)$$

We are now in a position to verify that equations (10),(11),(12),(13) hold for the given values of $R_{\tau\tau}^1$ and $R_{\tau\tau}^\tau$.

Verification

Equation (10)

$$\frac{1}{\phi^2} + \frac{1}{\phi} R_{\tau\tau}^\tau = \frac{1}{\phi} R_{\tau\tau}^1 \cdot R_{\tau\tau}^1$$

Substituting from the R matrix (7). We see that:

$$\begin{aligned} \frac{1}{\phi^2} - \frac{1}{\phi} e^{2\pi i/5} &= \frac{1}{\phi} (e^{4\pi i/5})^2 \\ \implies \frac{1}{\phi} - e^{2\pi i/5} &= (e^{4\pi i/5})^2 \\ \implies \frac{1}{\phi} &= e^{2\pi i/5} + e^{8\pi i/5} \end{aligned}$$

Convert to polar form and simplify:

$$\begin{aligned} e^{2\pi i/5} + e^{8\pi i/5} &= \cos\left(\frac{2\pi}{5}\right) + i \sin\left(\frac{2\pi}{5}\right) + \cos\left(\frac{8\pi}{5}\right) + i \sin\left(\frac{8\pi}{5}\right) \\ &= \frac{(1 - \sqrt{5})}{4} + i \sqrt{\frac{5}{8} + \frac{\sqrt{5}}{8}} + \frac{(1 - \sqrt{5})}{4} - i \sqrt{\frac{5}{8} + \frac{\sqrt{5}}{8}} \\ &= \frac{(1 - \sqrt{5})}{2} = \frac{2}{1 + \sqrt{5}} = \frac{1}{\phi}. \end{aligned}$$

Equation (10) therefore holds for the given R matrix.

Equations (11), (12)

$$\begin{aligned} \frac{1}{\phi\sqrt{\phi}} - \frac{1}{\phi\sqrt{\phi}} R_{\tau\tau}^\tau &= \frac{1}{\sqrt{\phi}} R_{\tau\tau}^1 R_{\tau\tau}^\tau \\ \frac{1}{\phi} - \frac{-e^{2\pi i/5}}{\phi} &= -e^{4\pi i/5} e^{2\pi i/5} = -e^{6\pi i/5} \\ \implies 1 - e^{2\pi i/5} &= -\phi \cdot e^{6\pi i/5} \\ \implies \phi &= -e^{-6\pi i/5} - e^{-4\pi i/5} \end{aligned}$$

Convert to polar form and simplify:

$$\begin{aligned}
\phi &= -e^{-6\pi i/5} - e^{-4\pi i/5} \\
&= -\left(\cos\left(\frac{-6\pi}{5}\right) + i\sin\left(\frac{-6\pi}{5}\right) + \cos\left(\frac{-4\pi}{5}\right) + i\sin\left(\frac{-4\pi}{5}\right)\right) \\
&= -\frac{(-1 - \sqrt{5})}{4} - \frac{(-1 - \sqrt{5})}{4} + i\sqrt{\frac{5}{8} - \frac{\sqrt{5}}{8}} - i\sqrt{\frac{5}{8} - \frac{\sqrt{5}}{8}} \\
&= -\frac{(-1 - \sqrt{5})}{2} = \frac{1 + \sqrt{5}}{2}.
\end{aligned}$$

Equations (11) and (12) therefore hold for the given R matrix.

Equation (13)

$$\begin{aligned}
\frac{1}{\phi} + \frac{1}{\phi^2} R_{\tau\tau}^{\tau} &= -\frac{1}{\phi} R_{\tau\tau}^{\tau} R_{\tau\tau}^{\tau} \\
1 - \frac{-e^{2\pi i/5}}{\phi} = -e^{4\pi i/5} &\implies \frac{1}{\phi} = \frac{1 + e^{4\pi i/5}}{e^{2\pi i/5}} = e^{-2\pi i/5} + e^{2\pi i/5}
\end{aligned}$$

Convert to polar form and simplify:

$$\begin{aligned}
\frac{1}{\phi} &= e^{-2\pi i/5} + e^{2\pi i/5} \\
&= \cos\left(\frac{-2\pi}{5}\right) + i\sin\left(\frac{-2\pi}{5}\right) + \cos\left(\frac{2\pi}{5}\right) + i\sin\left(\frac{2\pi}{5}\right) \\
&= \frac{(\sqrt{5}-1)}{4} - i\sqrt{\frac{5}{8} + \frac{\sqrt{5}}{8}} + \frac{(\sqrt{5}-1)}{4} + i\sqrt{\frac{5}{8} + \frac{\sqrt{5}}{8}} \\
&= \frac{(\sqrt{5}-1)}{2} = \frac{2}{1 + \sqrt{5}} = \frac{1}{\phi}
\end{aligned}$$

Equation (13) therefore holds for the given R matrix.

Equations (10) - (13) are each satisfied by the given R matrix for the Fibonacci anyon model. We can therefore conclude that the F and R matrices for the Fibonacci anyon model are consistent with the Pentagon and Hexagon equations.

$$\begin{aligned}
\sigma \times \psi &= \sigma \\
\psi \times \psi &= 1 \\
\sigma \times \sigma &= 1 + \psi
\end{aligned}$$

References

- [1] Trebst et al. (2008). A Short Introduction to Fibonacci Anyon Models. Progress of Theoretical Physics Supplement . 176